Talk to be given at the minisymposium on "Computer Algebra" of the SIMAI 98-4th National Congress of the Italian Society for Applied and Industrial Mathematics, Giardini Naxos, Sicily 1-5 June 1998

Problems in Computing Properties of Symmetric Functions and Lie Groups

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> And yet the mystery of mysteries is to view machines making machines; a spectacle that fills the mind with curious, and even awful, speculation.

— Benjamin Disraeli: Coningsby (1844)

What are S-functions?

Suppose $(\mu) = (\mu_1, \mu_2, \dots, \mu_m)$ is a partition of an integer into integer parts μ_i then we can associate with it a *monomial*

$$\mathbf{x}^{\mu} = x_1^{\mu_1} x_2^{\mu_2} \dots x_m^{\mu_m} \tag{1}$$

Example:- $\mathbf{x}^{30126} = x_1^3 x_2^0 x_3^1 x_4^2 x_5^6 \equiv x_1^3 x_3^1 x_4^2 x_5^6$ Consider a *tableau* T of shape λ then define

$$\mathbf{x}^T = \prod_{(i,j)\in\lambda} x_{T_{i,j}} = \mathbf{x}^\mu \tag{2}$$

Example:- if

then

$$\mathbf{x}^T = x_1^3 x_2^2 x_3^2 x_5$$

A tableau T of shape λ is *semi-standard* if the integers appearing in rows are weakly increasing and strongly increasing down columns. The Schur-function (S-function) is defined by

$$s_{\lambda}(\mathbf{x}) = \sum_{T} \mathbf{x}^{T} \tag{3}$$

Example:- associated with $s_{21}(x_1, x_2, x_3)$ are the eight tableaux

and hence

$$s_{\lambda}(\mathbf{x}) = x_1^2 x_2 + x_1^2 x_3 + x_2^2 x_3 + x_1 x_3^2 + x_2 x_3^2 + x_1 x_2^2 + 2x_1 x_2 x_3$$
(4)

Placing no limits on the number of variables \mathbf{x} we can write

$$s_{21}(\mathbf{x}) = \sum_{i,j} x_i x_j + 2 \sum_{i,j,r} x_i x_j x_r$$
(5)

where the summations are carried out over all distinct permutations of the indices. Frequently we will designate an S-function $s_{\lambda}(\mathbf{x})$ by enclosing the partition (λ) in curly brackets $\{\lambda\}$ and leave the number of variables unspecified.

S-functions are symmetric functions, thus their products and powers may be resolved into sums of S-functions.

What can you do with S-functions?

1. Outer Products

$$\{\mu\} \cdot \{\nu\} = \sum_{\lambda} c^{\lambda}_{\mu\nu} \{\lambda\}$$
(1)

where the $c^{\lambda}_{\mu\nu}$ are non-negative integers known as the Littlewood-Richardson coefficients and the weights, ω_{λ} , are constrained by $\omega_{\lambda} = \omega_{\mu} + \omega_{\nu}$

2. Skews

$$\{\lambda/\mu\} = \sum_{\nu} c^{\lambda}_{\mu\nu}\{\lambda\}$$
(2)

The weights, ω_{ν} , are constrained by $\omega_{\nu} = \omega_{\lambda} - \omega_{\mu}$

3. Plethysms

$$\{\lambda\} \otimes \{\mu\} = \sum_{\nu} g^{\nu}_{\lambda\mu} \{\nu\}$$
(3)

where the $g^{\nu}_{\lambda\mu}$ are non-negative integers and the weights, ω_{ν} , are constrained by $\omega_{\nu} = \omega_{\lambda} \times \omega_{\mu}$

4. Inner Products

$$\{\mu\} * \{\nu\} = \sum_{\lambda} c^{\lambda}_{\mu\nu} \{\lambda\}$$
(4)

where the $c_{\mu\nu}^{\lambda}$ are non-negative integers and the weights of the partitions are constrained by $\omega_{\mu} = \omega_{\nu} = \omega_{\lambda} = n$.

5. Inner Plethysms

$$\{\mu\} \circ \{\nu\} = \sum_{\lambda} c^{\lambda}_{\mu\nu} \{\lambda\}$$
 (5)

where the $c_{\mu\nu}^{\lambda}$ are non-negative integers and the weights of the partitions are constrained by $\omega_{\mu} = \omega_{\lambda} = n$ and $\omega_{\nu} \ge 0$. Examples with SCHUR

```
SFN>
021,32
    \{53\} + \{521\} + \{4^2\} + 2\{431\}
  + \{42^2\} + \{421^2\} + \{3^22\} + \{3^21^2\}
  + \{32^2 1\}
SFN>
sk321,21
      \{3\} + 2\{21\} + \{1^3\}
SFN>
pl21,3
      \{63\} + \{531\} + \{52^2\} + \{521^2\}
    + \{4^2 1\} + \{432\} + \{431^2\} + 2\{42^2 1\}
    + \{421^3\} + \{41^5\} + \{3^3\} + \{3^2, 21\}
    + \{3^2 1^3\} + \{32^3\} + \{32^2 1^2\}
SFN>
i32,2111
      \{32\} + \{31^2\} + \{2^2, 1\} + \{21^3\}
SFN>
i_p121
      <21> + <2> + <1^2 > + <1>
SFN>
```

Some infinite series of S-functions

$$\begin{split} L &= \sum_{m=0}^{\infty} (-1)^m \{1^m\} \quad M = \sum_{m=0}^{\infty} \{m\} \\ P &= \sum_{m=0}^{\infty} (-1)^m \{m\} \quad Q = \sum_{m=0}^{\infty} \{1^m\} \\ B &= \sum_{\beta} \{\beta\} \quad D = \sum_{\delta} \{\delta\} \end{split}$$

where the *m* are integers, the partitions (δ) are all partitions having only even parts while the partitions (β) are conjugates of the (δ) . *L* and *M* are inverses of one another as are *P* and *Q*.

SCHUR can compute these series, and many others, up to a user determined limit.

Examples:-

New S-function identities

Infinite S-function series play a key role in practical calculations for both compact and non-compact Lie groups. SCHUR gave evidence leading to a number of conjectures involving plethysms of certain infinite S-functions.

$$M_{+} = \sum_{m=0}^{\infty} \{2m\} \quad M_{-} = \sum_{m=0}^{\infty} \{2m+1\}$$
$$L_{+} = \sum_{m=0}^{\infty} \{1^{2m}\} \quad L_{-} = \sum_{m=0}^{\infty} \{1^{2m+1}\}$$
(1)

$$A_{\pm} = \{1^2\} \otimes L_{\pm} \quad B_{\pm} = \{1^2\} \otimes M_{\pm} \\ C_{\pm} = \{2\} \otimes L_{\pm} \quad D_{\pm} = \{2\} \otimes M_{\pm}$$
(2)

Let
$$Z_{\pm} = \{A_{\pm}, B_{\pm}, C_{\pm}, D_{\pm}\}$$
 then
 $Z_{+} \otimes \{1^{2}\} = Z_{-} \otimes \{2\}$
 $Z_{+} \otimes \{21^{2}\} = Z_{-} \otimes \{31\}$
(3)

- M Yang and B G Wybourne, J. Phys. A: Math. Gen. 19 3513 (1986)
- R C King, B G Wybourne and M Yang, J. Phys. A: Math. Gen. 22 4519 (1989)
- 3. K Grudzinski and B G Wybourne, J. Phys. A: Math. Gen. **29** 6631 (1996)

A tensor product in SO(10)

$$[\lambda] \times [\mu] = \sum_{\zeta} [\lambda/\zeta \cdot \mu/\zeta]$$
(1)

Example:-

$$[1^{3}] \times [2^{3}] = \sum_{\zeta} [1^{3}/\zeta \cdot 2^{3}/\zeta]$$
$$[1^{3} \cdot 1^{3}] + [1^{2} \cdot 2^{2}1] + [1 \cdot 21^{2}] + [0 \cdot 1^{3}]$$

$$[1^{3} \cdot 2^{3}] = [3^{3}] + [3^{2}21] + [32^{2}1^{2}] + [2^{3}1^{3}]$$

$$[1^{2} \cdot 2^{2}1] = [3^{2}1] + [32^{2}] + [321^{2}] + [2^{3}1] + [2^{2}1^{3}]$$

$$[1 \cdot 21^{2}] = [31^{2}] + [2^{2}1] + [21^{3}]$$

$$[0 \cdot 1^{3}] = [1^{3}]$$

$$[32^{2}1^{2}] \equiv [32^{2}1^{2}]_{+} + [32^{2}1^{2}]_{-}$$
$$[2^{3}1^{3}] \equiv [2^{3}1]$$
$$[2^{2}1^{3}] \equiv [2^{2}1^{3}]_{+} + [2^{2}1^{3}]_{-}$$

$$\begin{split} [1^3] \times [2^3] &= [3^3] + [3^2 21] + [3^2 1] + [32^2 1^2] + \\ &+ [32^2 1^2] - + [32^2] + [321^2] + [31^2] \\ &+ 2[2^3 1] + [2^2 1^3] + + [2^2 1^3] - + [2^2 1] \\ &+ [21^3] + [1^3] \end{split}$$

Calculation by SCHUR

```
gr so10
Group is SO(10)
REP>
p111,222
    [3^3] + [3^2 21] + [3^2 1] + [32^2 1^2]+
+ [32^2 1^2] - + [32^2] + [321^2] + [31^2]
+ 2[2^3 1] + [2^2 1^3] + [2^2 1^3] -
+ [2^2 1] + [21^3] + [1^3]
REP>
dim last
dimension=495000
```

- 1. Time taken to compute the result by hand < 2 minutes.
- 2. Time taken by SCHUR on a Pentium instantaneous.
- 3. Time reported in the literature on a VAX4000 5 hours CPU.

The product $[6^3] \times [9^3]$ is of dimension 92,908,920,088,670,400. SCHUR resolved the product in 40 minutes on a SUN IPX. Undoubtedly the Vax4000 programme would take a time that would dwarf the age of the universe!

S-function series and branching rules

 $U(n) \downarrow U(n-1)$ $\{\lambda\} \downarrow \{\lambda/M\}$ $U(n) \downarrow O(n)$ $\{\lambda\} \downarrow [\lambda/D]$ $U(2n) \downarrow Sp(2n)$ $\{\lambda\} \downarrow \langle \lambda/B \rangle$

$$Sp(2n, R) \downarrow U(n)$$

$$\langle \frac{k}{2}(\lambda) \rangle \downarrow \varepsilon^{\frac{k}{2}} \cdot \{\{\lambda_s\}_N^k \cdot D_N\}_N \quad N = \min(k, n)$$

$$SO^*(2n) \downarrow U(n)$$

$$[k(\lambda)] \downarrow \varepsilon^k \cdot \{\{\lambda_s\}_N^{2k} \cdot B_N\}_N \quad N = \min(2k, n)$$

Examples of branching rules with SCHUR

```
DP>
gr u4
Group is U(4)
DP>
br1,4gr1[321]
Group is O(4)
       [31] + [2^2] + [2] + [2] + [1^2]
DP>
gr u4
Group is U(4)
DP>
br2,4gr1[321]
Group is Sp(4)
      <31> + <2^2 > + <2> + <1^2 >
DP>
gr spr6
Group is Sp(6,R)
DP>
br36,6gr1[2;21]
Group is U(3)
 \{432\} + \{4^2 3\} + \{53^2 \} + \{542\} + \{54^2 \}
 + \{5^2 3\} + \{632\} + 2\{643\} + \{652\} + 2\{654\}
 + \{6^2 3\} + \{6^2 5\} + \{73^2 \} + \{742\}
 + \{74^2\} + 2\{753\} + \{75^2\} + \{762\} + \dots
```

S-functions and tensor products

$$\begin{split} U(n): \{\mu\} \times \{\nu\} &= \sum_{\lambda} C_{\mu\nu}^{\lambda} \{\lambda\} \\ O(n): [\mu] \times [\nu] &= \sum_{\zeta} [\mu/\zeta \cdot \nu/\zeta] \\ Sp(2n): \langle\mu\rangle \times \langle\nu\rangle &= \sum_{\zeta} \langle\mu/\zeta \cdot \nu/\zeta\rangle \\ Sp(2n, R): \langle\frac{k}{2}(\mu)\rangle \times \langle\frac{\ell}{2}(\nu)\rangle &= \langle\frac{k+\ell}{2}(\{\mu_s\}^k \cdot \{\nu_s\}^\ell \cdot D)_{k+\ell,n}\rangle \\ SO^*(2n): [k(\mu)] \times [\ell(\nu)] &= [k+\ell(\{\mu_s\}^{2k} \cdot \{\nu_s\}^{2\ell} \cdot B)_{k+\ell,n}] \end{split}$$

```
REP>
gr sp6
Group is Sp(6)
REP>
p21,31
  <52> + <51^2 > + <5> + <43> + 2<421>
 + 3<41> + <3^2 1> + <32^2 > + 3<32>
 + 3<31^2 > + 2<3> + 2<2^2 1> + 3<21>
 + <1^{3} > + <1>
REP>
gr so8
Group is SO(8)
REP>
p s;0+,21
 [s;21] + + [s;2] - + [s;1^2] - + [s;1] +
REP>
gr spr6
Group is Sp(6,R)
REP>
p1;0,2;1
  <3;(1)> + <3;(1^3)> + 2<3;(21)>
 + \langle 3; (2^2 1) \rangle + \langle 3; (3) \rangle + \langle 3; (31^2) \rangle
 + 2<3; (32)> + <3; (3<sup>2</sup> 1)> + 2<3; (41)>
 + \langle 3; (421) \rangle + 2\langle 3; (43) \rangle + \ldots
```

Algebraic approaches to the genetic code

Hornos and Hornos¹ investigated those simple Lie algebras having at least one representation of dimension 64, the number 64 corresponding to the $4 \times 4 \times 4$ possible codons, each involving four bases arranged in triplets, to code the 20 amino acids.

The groups Sp(6) and G(2) were found^{1,2} to be of particular interest. SCHUR has been able to determine the various possible group-subgroup decompositions and the eigenvalues of the Casimir operators used to describe the possible symmetry breakings.

In addition SCHUR was used to establish the complete set of 64-dimensional representations for the symmetric and alternating groups.

The complete set of 64–dimensional irreducible representations for the groups S(n) and A(n)

$ \begin{bmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 &$	$ \begin{array}{ c c c c c c c c c c c c c c c c c c c$	$\{521\}\ [521]\ \{64\ 1\}\ [64\ 1]$	$\{321^3\}$ $\{21^{63}\}$	$S(13) \\ S(14) \\ A(14) \\ A(15)$	$ \begin{array}{c} \{\Delta\} \\ \{\Delta_{\pm}\} \\ [64 \ 1] \\ [\Delta_{\pm}] \end{array} $
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- J E M Hornos and Y M M Hornos, Phys. Rev. Lett. 71 4401 (1993)
- M Forger, Y M M Hornos and J E M Hornos, Phys. Rev. E56 7078 (1997)
- 3. R D Kent, M Schlesinger and B G Wybourne, Can. J. Phys. (In Press)

Problems in symplectic models of nuclei, quantum dots and many-electron states often involve the symmetric group S_n . Applications require the resolution of symmetrised powers of tensor representations of S_n . These are required in determining branching coefficients. The coefficients involve *inner plethysms*. Of particular interest is the representation

$$\{n-1,1\} \equiv \langle 1 \rangle \tag{1}$$

and the inner plethysms

$$\langle 1 \rangle \otimes \{\lambda\} = \sum_{\rho} c_{\lambda}^{\rho} \langle \rho \rangle \tag{2}$$

SCHUR has computed the complete resolution of the plethysms $\langle 1 \rangle \otimes \{n\}$ for n = 1, ..., 20. Applications often require the value of single coefficients in very large plethysms. Here generating methods can be used. Thus with MAPLE it was possible to show that in S_{40}

$$\{39,1\} \otimes \{30\ 4321\} \supset 309,727,790,880\{31\ 3^22\}$$

A calculation quite beyond SCHUR.

- T Scharf, J-Y Thibon and B G Wybourne, J. Phys. A: Math. Gen. 26 7461 (1993)
- 2. T Scharf, J-Y Thibon and B G Wybourne, J. Phys. A: Math. Gen. **30** 6963 (1997)

The Vandermonde determinant and the quantum Hall effect

The Vandermonde alternating function in N variables is defined as

$$V(z_1, \dots, z_N) = \prod_{i < j}^{N} (z_i - z_j)$$
(1)

Any even power, V^{2m} , is necessarily a *symmetric* function and hence expandable into a set of symmetric functions such as the Schur functions

$$s_{\lambda}(z_1, \dots, z_N) = \{\lambda\} = \{\lambda_1, \dots, \lambda_p\}$$
(2)

which in this case are indexed by partitions of the integer

$$n = mN(N-1) \tag{3}$$

We need the expansion coefficients c^{λ} for

$$V^{2m} = \sum_{\lambda \vdash n} c^{\lambda} s_{\lambda} \tag{4}$$

where the c^{λ} are signed integers and are precisely the same integers that arise in the expansion of the Laughlin wavefunction, used in the quantum Hall effect, as a linear combination of Slater determinants.

This is a COMBINATORIALLY EXPLOSIVE problem!

 T Scharf, J-Y Thibon and B G Wybourne, J. Phys. A: Math. Gen. 27 4211 (1994).

N	$N_{tableaux}$	$N_{tableaux}^{conjectured^*}$	N_{coeff}
1	1	1	1
2	2	2	4
3	5	5	28
4	16	16	292
5	59	59	$4,\!102$
6	247	247	$73,\!444$
7	$1,\!111$	$1,\!111$	$1,\!605,\!838$
8	$5,\!294$	$5,\!302$	41,603,200
9	$26,\!310$	$26,\!376$	

* "The above reasoning does not however insure that this is exactly the total number of tableaux in the expansion of V^{2m} in characters as some coefficients might still vanish. However experience up to N = 5 seems to indicate that these accidents do not happen" P. Di Francesco, M. Gaudin, C. Itzykson and F. Lesage. (SphT/93-125)

Invariants formed from the Riemann tensor

The master object for enumerating Riemann scalars is

$$\mathcal{G} = \sum_{m=1}^{\infty} (t^2 \{2^2\} + t^3 \{32\} + t^4 \{42\} + \ldots + t^p \{p, 2\} + \ldots)^m \quad (1)$$

- 1. There is a Riemann scalar for every S-function $\{\lambda\}$ arising in (1) whose partition label $\lambda = \lambda_1, \lambda_2, \ldots, \lambda_p$ involves only *even* parts.
- 2. The evaluation of the Riemann scalars of order n involves the resolution of all plethysms and outer S-function products associated with t^n where n is necessarily even.

Order <i>n</i>	Number of Riemann Scalars	
2	1	
4	4	
6	17	
8	92	
10	668	
12	6,721	
14	89,137	

- 1. S A Fulling, R C King, B G Wybourne and C J Cummins, Class. Quantum Grav. **9** 1151 (1992)
- B G Wybourne and J Meller, J. Phys. A: Math. Gen. 25 5999 (1992)

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Research supported by Polish KBN Grants

Questions?

The only questions worth asking are the unanswerable ones — John Ciardi Saturday Review-World (1973)

For every complex question there is a simple answer — and it's wrong.

— H. L. Mencken